# AN ATTEMPT AT AN ALGEBRAIC THEORY OF CRYSTAL STRUCTURE, PART 3.

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### ABSTRACT

The aim of this work is to found a theory of crystal structure on the notions not of length and angle, but of number, the fundamental phenomenon involved being the patterns produced by X-rays falling on matter in the solid, or crystalline state. It is shown that some alloys of transition metals are a source of patterns complex and varied enough to enable gaps in their enumeration to be exposed by the rudiments of such a theory (Aboav, 1997, 1998).

Key words: crystal structure, algebraic theory, space-filling polyhedra.

#### INTRODUCTION

In Parts 1 & 2 (Aboav, 1997, 1998) it was shown how the topological properties of some inclusion hydrates of the 'dodecahedral series' could be classified in terms of the solutions of a simple Diophantine equation, and the result used to predict the existence of further compounds of the same type. These properties are summarized in Table 1, which shows for five of the hydrates the number of 16–, 15–, 14–, and 12–hedra (cols.2 through 5);  $N_0$ , the number of vertices (col.6);  $N_3$ , the number of polyhedra (col.7) in a unit cell of their crystal structure; and (col.8) the number  $29N_3 - 5N_0$ .  $N_0$  and  $N_3$  made relatively prime by the removal of their common factors.

Table 1. Topological properties of some inclusion hydrates.

Tubic	1. TOP	ological	010	per ties or	301116	IIICIUS	ion nyurates.		
hydrate	16	15	14	12	N <sub>0</sub>	N <sub>3</sub>	$(29N_3 - 5N_0)$	no	n <sub>3</sub>
chlorine			6	2	46	8	2	23	4
chlorofrom	8			16	136	24	16	17	3
alkyl onium sal	t	4	4	6	80	14	6	40	7
bromine tetr.		4	16	10	172	30	10	86	15
bromine ortho.	4	4	4	14	148	26	14	74	13

The table shows that for each of the five unit cells listed:

(i) the number of 12-hedra is equal to  $29N_3 - 5N_0$ , and

(ii) the numbers of 16-, 15-, 14-, & 12-hedra are multiples, or sums of multiples of their combinations  $1 \mid 0 \mid 0 \mid 2$ ,  $0 \mid 2 \mid 2 \mid 3$ , and  $0 \mid 0 \mid 3 \mid 1$ , thus:

chlorine 2x(0|0|3|1), chloroform 8x(1|0|0|2), alkyl onium salt 2x(0|2|2|3), bromine (tetr.) 4x(0|0|3|1) + 2x(0|2|2|3), & bromine (orth.) 2x(0|2|2|3) + 4x(1|0|0|2).

These two findings are, however, not independent; for, if the multiples of 1|0|0|2, 0|2|2|3, and 0|0|3|1 are integers i, j, and k, respectively, the number  ${}_{12}N_3$  of 12-hedra in a unit cell is given by

$$_{12}N_3 = 2i + 3j + k;$$
 (1)

and, since 16-, 15-, 14-, and 12-hedra possess 28, 26, 24, and 20 vertices --- or, if they are so placed that each vertex is shared by four contiguous polyhedra, 7,  $6\frac{1}{2}$ , 6, and 5 vertices, respectively --- the number  $N_0$  of vertices of such a cell is given by

$$N_0 = 17i + 40j + 23k$$
 (2)

and the number N<sub>3</sub> of its constituent polyhedra by

$$N_3 = 3i + 7j + 4k$$
; (3)

from which it follows that

$$29N_3 - 5N_0 = 2i + 3j + k$$

$$= {}_{12}N_3.$$
(4)

This shows that, if  $N_0$  and  $N_3$  are known, the number of 12-hedra in a unit cell and consequently the *total number*  $(N_3 - {}_{12}N_3)$  of its remaining polyhedra are also known. But the *partition* of the latter into 16-, 15-, and 14-hedra is not deducible from  $N_0$  and  $N_3$ , since the operation (1 | -2 | 1 | 0) --- i.e. the replacement of two 15-hedra by a 16-hedron and a 14-hedron --- leaves  $N_0$  and  $N_3$  invariant and, since

$$(1 | -2 | 1 | 0) \equiv (1 | 0 | 0 | 2) - (0 | 2 | 2 | 3) + (0 | 0 | 3 | 1),$$
 (5)

allows the numbers of different polyhedra, after the operation, to be expressible as multiples, or sums of multiples of the same combinations as before.

The above rule (4) does not, however, hold generally. It is a feature of 4-connected honeycombs whose polyhedral cells have pentagonal and non-adjacent hexagonal faces only; and many inclusion hydrates in particular (Jeffrey, 1982) have structures that do not obey it.

But there are compounds, among them alloys of the transition metals, whose structure, though differing chemically from that of the hydrates, is nevertheless governed by that rule; and the object of the investigation now to be described is to show how those compounds, too, can be classified and their existence predicted by means of the above-mentioned Diophantine equation.

## THE DIOPHANTINE EQUATION $2^ax = 3^by + z$ .

It has already been shown (Aboav, 1998) that the quantities  $N_0$  and  $N_3$ , and hence  $n_0$  and  $n_3$ , of Table 1 are proportional to the roots of a Diophantine equation:

$$2^{a}x = 3^{b}y + z \tag{6}$$

where a and b are positive integers, and x, y, and z odd primes or unity, with z positive or negative. To obtain agreement with experiment certain restrictions on the values of x, y, and z had to be made, viz:

(i) 
$$x, y < 100$$
,

(ii) 
$$z/2^a x < 1/100$$
;

but the experimental facts now to be investigated suggest a change should be made both in the composition and in the limits of x, y, z.

As regards their composition each of these numbers was assumed to be prime (Aboav, 1997); but it will now be assumed that:

- (a) z is composite,
- (b) one of the two numbers x y is odd prime, while the other may be odd prime or the product of two odd primes, of which one may not exceed 7.

This gives rise to five equations to be solved for x and y:

$$2^{a}x = 3^{b}y + z$$

$$2^{a}5.p = 3^{b}y + z$$

$$2^{a}x = 3^{b}5.q + z$$
(7)

$$2^{a}7.p = 3^{b}y + z$$

$$2^{a}x = 3^{b}7.q + z$$
(8)

where p and q are odd primes. So that they may be more easily recognized, composite numbers are shown underlined.

As regards the limits imposed on x and y it will now be assumed that:

(iii) 
$$x < 240$$
;  $y < 60$ .

Since z is assumed composite there is no longer need to refer to its value explicitly; so that restriction (ii) may be expressed as  $1 - 3^by/2^ax < 1/100$ , or as

(iv) 
$$1 - 3^6 n_3 / 2^7 n_0 < 1/100$$

(Aboav, 1997, 1998); and the solutions of Eq.(6) which were formerly denoted by  $[a\ b]\{x\ y\ z\}$  may be written simply  $[a\ b]\{x\ y\}$ .

The fraction  $(1 - 3^6 n_3/2^7 n_0)$  will henceforth be referred to as the *comma*, and denoted by the Greek letter  $\kappa$ :

$$K = 1 - 3^6 n_3 / 2^7 n_0 \tag{9}$$

so that restriction (iv) can be written

(v) 
$$\kappa < 1/100$$
.

For greater clarity the comma may sometimes here be expressed as the nearest vulgar fraction with unity as numerator, as in Table 5A below.

24 solutions to Eq.(6) subject to these new conditions are shown in Table 2, which contains many of the solutions listed in Part 2 (Aboav 1998) as well as some further ones. In some cases more than one solution is shown for a given [a b].

Table	2. Solutions	of Eq.(6) subject	to restrictions	s (iii) and (iv
	b = 4	b = 5	b = 6	b = 7
a 	{ x y }	{ x y }	{ x y }	{ x y }
5	{43 17}	{53 7}	<b>{23</b> 1 <b>}</b>	
6	{29 23}	{19 5}	{149 13}	{103 3}
7	{7 11}  	{59 31} {89 47} {101 53}	{97 17}	
8	{13 41} 	{41 43} 	{37 13} {151 53}	{43 5} 
9	{ 3 19}	{11 23}	{67 47}	{47 11}
10			<b>{ 5 7</b> }	

Table 2. Solutions of Eq.(6) subject to restrictions (iii) and (iv)

13 solutions of Eq.(7) and 7 solutions of Eq.(8) are shown in Tables 3 & 4, respectively. In Tables (2), (3) & (4) numbers that describe a known crystal structure  $(v.\ infra)$  are shown in **bold type**.

These tables, to which more solutions subject to the same restrictions can be added, are not meant to be complete. The 44 solutions here chosen are gathered together in Table 5A, in which are listed, after a serial number (col.1) a solution (col.2) and its corresponding values of  $n_0$  and  $n_3$  (cols.3 & 4). Column 5 lists the comma,  $\kappa$ ; cols.6 & 7, the numbers  $60(1 + n_3/n_0)$  and  $2(1 + n_0/n_3)$  to 3 decimal places; and col.9, the integer  $29n_3-5n_0$ .

Table 3. Solutions of Eq.(7) subject to restrictions (iii) and (iv)								
	b = 5	b = 6	b = 7	b = 8	b = 9			
a	{ x y}	{ x y}	{ x y}	{ x y	} { x y}			
5			{ <u>5.41</u> 3}					
6								
7	$   \left\{ \frac{5.11}{67} \right. 29 $	} { <u>5.47</u> 41} }		{ <u>5.31</u> 3	}			
8	{ <u>5.7</u> 37}	{71 <u>5.5</u> } {157 <u>5.11</u> }	{ <u>5.29</u> 17}					
9	{31 <u>5.13</u>	}	{107 <u>5.5</u> }		{ <u>5.23</u> 3}			
10	{13 <u>5.11</u>	_}						
Table	4. Solutio	ns of Fa.(8)	subject to re	estrictions	(iii) and (iv)			
Table		ns of Eq.(8) b = 5			(iii) and (iv) b = 8			
	b = 4		b = 6	b = 7	b = 8			
	b = 4	b = 5	b = 6 { x y}	b = 7	b = 8			
а	b = 4	b = 5 { x y}	b = 6 { x y}	b = 7	b = 8			
а 3	b = 4	b = 5 { x y}	b = 6 { x y}	b = 7	b = 8			
3 4	b = 4	$b = 5$ $\{ x \ y \}$ $\{ 7.13 \ 3 \}$	b = 6 { x y}	b = 7	b = 8			
3 4 5	b = 4	$b = 5$ $\{ x \ y \}$ $\{ 7.13 \ 3 \}$	b = 6 { x y} +	b = 7	b = 8			
3 4 5	b = 4 { x y }  {31 7.7}	$b = 5$ $\{ x \ y \}$ $\{ 7.13 \ 3 \}$	b = 6 { x y} +   7.31 19}	b = 7 { x y}	b = 8			
3 4 5 6	b = 4 { x y }  {31 7.7}	b = 5  { x y}  { <u>7.13</u> 3}  {	b = 6 { x y} +	b = 7 { x y}	b = 8 { x y} {7.11 3}			

THE CRYSTAL: THE PICTURES OF HAÜY (1784), AND OF FEDOROV (1890).

The numbers listed in Table 5A, which are solutions of a Diophantine equation or numbers derived from them, are dimensionless and hence <code>scale-free</code>. It is proposed next to compare them with some other numbers which, since they depend on a physical process and necessitate measurements of length and angle, are <code>scale-dependent</code>.

Table 5A. Solutions to Eqs (6),(7), & (8), and some related numbers. No.  $(ab)\{xy\} n_0 n_3$  $60(1+n_3/n_0) 2(1+n_0/n_3) 29n_3-5n_0$ ĸ [5 6]{23 1} 23 4 1/104 70.435 13.500 1 1 (7 5){67 <u>5.7</u>} 201 35 1/121 70.448 13 - 486 10 2 13.481 3 [7 8]{5.31 3} 155 27 1/126 70.452 8 7 4 [9 5]{11 23} 132 23 1/131 70 - 455 13 - 478 5 5 [8 7]{43 5} 86 15 1/151  $70 \cdot 465$  $13 \cdot 466$ \*\* \*\* 1/151 70.465 13.466 5 1/157  $(7 \ 6)\{\underline{5.47} \ 41\} \ 235 \ 41$ 70.468 13 - 463 14 6 70.470 7 149 26 1/162 13.462 a [6 6]{149 13} [7 4]{7 11} 63 11 1/179 70 - 476 13 - 455 4 [9 5]{31 5.13} 372 65 1/206 70.484 13.446 25 (6 7){103 3} 103 18 1/213 70.485 13.444 7 1.0 1/223 11 [8 5]{41 43} 246 43 70.488 13.442 17 12 (9 7){7.19 31} 532 93 1/228 70.489 13 - 441 37 13 [7 5]{101 53} 303 53 1/264 70.495 13.434 22 14 [10 6]{5 7} 40 7 1/301 70.500 13 - 428 3 1/301 70.500 3 13 - 428 . \*\* 1/301 70.500 13.428 3 217 38 15 (6 6){7.31 19} 1/375 70.507 13 - 421 17 16 (7 5){59 31} 177 31 1/397 70.508 13 - 419 14 17 [8 6]{157 <u>5.11</u>} 314 55  $\textbf{70} \cdot \textbf{510}$ 1/414 13 - 418 25 18 [8 4]{13 41} 234 41 1/476 70.513 13 - 415 19 19  $(9\ 7)\{107\ \underline{5.5}\}$ 428 75 1/503 70.514 13 - 413 35 97 17 70.515 20 [7 6]{97 17} 1/540 13.412 8 1/680 70.519 13 [8 8]{7.11 3} 154 27 13.407 21 22 (9 6){67 47} 268 47 1/837 70.522 13 - 404 23 23 [6 5]{19 5} 57 10 1/1216 70.526 13.400 5 [8 6]{151 53} 302 53 1/2035 70.530 13.396 27 24 70.532 2.5 (9 7){47 11} 188 33 1/3438 13.394 17 26 [7 6]{131 23} 131 23 1/16,768 70.534 13.391 12 27 (5 7){5.41 3} 205 36 -1/6562 70.537 13.389 19 28 [7 4]{31 7.7} 279 49 -1/3968 70.538 13.388 26 74 13 13.385 7 29 (8 6){37 13} -1/1894 70.540 70.540 7 -1/1894 13.385 ,, -1/1894 70.540 13.385 7 . . -1/1894 70.540 13.385 7 387 68 70.543 30 (5 4){43 17} -1/1376 13.382 37  $70\cdot 546$ 31 [7 5]{5.11 29} 165 29 -1/1006 13.379 16 438 77 -1/813 70.548 13.377 43 32 [8 5]{73 7.11} 33 (3 5){7.13 3} 91 16 -1/728 70.549 13.375 9 70.552 29 34  $(8 7)\{\underline{5.29} 17\}$ 290 51 -1/629  $13 \cdot 373$ 70.556 35 [9 4]{3 19} 108 19 -1/512 13.368 1 1 392 69 -1/401 70.561 13.362 41 36 [10 7]{7.7 23} 37 [7 5]{89 47} 267 47 -1/393 70.562 13.362 2.8 [8 6]{71 <u>5.5</u>} 142 25 -1/371 70.563 13.360 15 38 70.565 39 [9 9]{5.23 3} 460 81 -1/348 13 - 358 49 40 (5 5) {53 7} 159 28 -1/339 70.566 13.357 17 70.571 41  $\{8\ 5\}\{5.7\ 37\}$ 210 37 -1/289 13.351 23 [6 4]{29 23} 261 46 -1/265 42 70.575 13.348 29 -1/251 70.577 35 43  $(10 5)\{13 5.11\} 312 55$ 13.345 44 [7 6]{17 3} 17 3 -1/198 70.588 13.333 2

Table 5B. Crystal structure of some 't.c.p.' alloys (Shoemaker & Shoemaker, 1986).

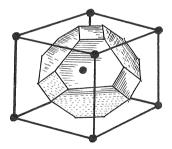
<u>1a</u>	ble 5B. Crysta	structure of sc	ome 't.c.p.' alloys (S	hoemaker	& Shoem	<u>aker, 1986).</u>
No		a(A) $b(A)$	c(Å) β(°) N	< 0 >	< CN >	_xcalc
1	Cr <sub>3</sub> Si	(4.564)	8	70.435	13.500	1 · 00
2						
3						
4						
5	$Cr_{46}Fe_{54}$	(8.800	4.544) 30	70.465	13 - 466	4 · 98
	H complex	(4·5 17·5	4.5) 30	70 - 465	13 · 466	4 · 98
6			82		13.46	
7	F complex		52		13.46	
8	J complex		22		13 · 45	
9						
10						
11						
12	Mn <sub>81.5</sub> Si <sub>18.5</sub>	5(16·992 28·63	4.656) 186	70 - 489	13 · 441	37 · 11
13	01 0 10 .	,				
14	Zr <sub>4</sub> Al <sub>3</sub>	(5.433	5 · 390) 7	70.500	13 - 428	3 · 00
		(9.070 16.983	4.752) 56	70.500	13 - 428	6 · 04
	MoNi		8 · 852) 56	70.500	13 - 428	6 · 04
15			·			
16						
17	Mn <sub>77</sub> Fe <sub>4</sub> Si <sub>19</sub>	(13.362 11.645	8 · 734 90 · 5)220	70.510	13 · 418	25 · 14
18	7, 4 13					
19						
20						
21						
22						
23						
24	Mo31Cr18Co51	1 (10.903	19 · 342) 159	70.530	13.396	26.93
25		-				
26						
27						
28						
29	Mo <sub>6</sub> Co <sub>7</sub>	(4.762	25 · 615) 39	70.540	13.385	6 · 96
	K7Cs6	(9.078	32.950) 26	70.540	13.385	6 · 96
	W <sub>6</sub> (Fe,Si) <sub>7</sub>	(9.283 7.81	7 4.755) 26	70.540	13.385	6 · 96
	Nb48Ni39Al1	13 (9.303 16.2	66 4.933) 52	70.540	13.385	6 · 96
30						
31						
32						
33						
34						
35	V <sub>41</sub> Ni <sub>36</sub> Si <sub>23</sub> (	13.462 23.381	8 · 940 100 · 3)228	70.556	13.369	10.94
36						
37						
38	V <sub>2</sub> (CoSi) <sub>3</sub>	(17 · 17 4 · 66	7 · 55 99 · 2) 50	70.563	13.360	14.92
39	Mg <sub>32</sub> (Zn,Al) <sub>3</sub>	39 (14·16)	162	70.565	13.358	48.75
40						
41	$Mn_{45}Co_{40}Si_1$	5 (15.42 12.	3 4 · 74) 74	70.571	13.351	22.90
42						
43	$Mg_{4}Zn_{7}$	(25.96 5.24	14.28 102.5) 110	70 · 577	$\textbf{13} \cdot \textbf{345}$	34.89
44	$MgZn_2(MgCu_2)$	(5.16(7.08)	8.50) 12(24)	70.588	13.333	2 · 00

The physical process referred to is the scattering of X-rays by matter in the crystalline state. The characteristic pattern of spots this produces on a photographic plate is thought of as originating in a microscopic structure possessed by the crystal. This postulated structure is pictured macroscopically in at least two ways. In the first way, associated with the name of Haüy (1784), the crystal is pictured as a regular array of particles, or 'atoms', located in a space of three dimensions: in the second, that of Fedorov (1890), the space is pictured as divided into contiguous, congruent polyhedral regions, or *stereohedra*. The two representations are equivalent, the one being the 'dual' of the other; but Fedorov's has for us the advantage that, in emphasizing a topological feature of crystal structure, it facilitates the eventual passage to an algebraic theory.

Fedorov's stereohedron is topologically equivalent to the *Voronoi polyhedron* (Okabe et al., 1992) generated by a point at its centre; and in some cases, as in the truncated octahedron of Fig.1, may also be congruent to it. Hence, where only a topological feature of a crystal is in question, Voronoi's figure can be used in place of Fedorov's. The number of faces of the Voronoi polyhedron is also known as the *coordination number*. CN. of the atom at its centre.

Figure 1 shows part of the structure of a body-centred cubic crystal pictured in each of the above two ways. The black dots represent the sites of Haüy's atoms; while the stereohedron of Fedorov, or Voronoi polyhedron, is shown as a truncated octahedron surrounding an atom at the centre of the cube.

Fig.1. Structure of body-centred cubic crystal.



In this case the stereohedra, like the polyhedral compartments of the crystal structures so far investigated (Aboav, 1997), form a 4-connected honeycomb; but in other structures, like for example the face-centred cubic lattice, they may be n-connected, with n > 4. That the algebraic theory now put forward may nevertheless be valid for these cases, too, may be seen from the following considerations.

Figure 2a shows a Voronoi diagram in two dimensions only: it is a periodic tessellation of 3-connected Voronoi polygons whose generating points, here represented as black dots, form a 'dual' tessellation of parallelograms. With the angle  $\beta$  of a parallelogram equal to 75°, as in the figure, the side AB of the Voronoi polygon is roughly 4mm long; and with  $\beta$  increased to 85° (Fig.2b) it diminishes to about 1mm.

As  $\beta$  approaches a rightangle there comes a stage where AB can no longer be distinguished from a point, and 4 instead of 3 sides of the Voronoi polygon, which now looks like a rectangle, appear to meet at the vertex A. (Fig.2c). The figure

may then be regarded either as a 4-connected, or as a 'degenerate' 3-connected tessellation of polygons. A degenerate figure often appears when its generating points are regularly spaced (Okabe et al., 1992), the concepts of 'symmetry' and 'degeneracy' going hand in hand.

Since the notion of degeneracy applies equally to honeycombs of polyhedra in a space of 3 dimensions, every such honeycomb can in principle be regarded as 4-connected (Okabe et al., 1992), thus enabling the proposed algebraic theory of crystal structure (Aboav, 1997), which invokes a geometrical property of 4-connected honeycombs, to be applied more generally.

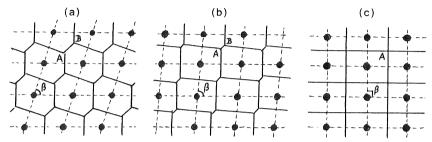


Fig.2. Periodic tessellations of Voronoi polygons in 2 dimensions.

We end this section by indicating briefly how the rudiments of such a theory may predict for some transition metals the topological structure of a number of alloys as yet unidentified.

### 'TETRAHEDRALLY CLOSE-PACKED' ALLOYS

The alloys of some transition metals, the so called 'tetrahedrally close-packed', or t.c.p. alloys, display a crystal structure of wide variety and complexity, which as we have seen may be described in terms either of the neighbouring atoms of a central atom, or of the space-filling packing of polyhedra. Shoemaker and Shoemaker (1986), used both these representations to express, for some 20 alloys of this type, metrical and topological features of their structure, of which some are reproduced in Table 5B

In col.1 of that table is shown an ordinal number, to help align its entries with those of Table 5A. The formula of each alloy is given in col.2, and its lattice constants in cols.3 through 6. N, the number of Voronoi polyhedra in its unit cell, is shown in col.7;  $<\Theta>$ , the average dihedral angle of the tetrahedral interstices formed by the atoms of the alloy, in col.8; and <CN>, the average coordination number of those atoms, in col.9. An approximate value,  $x_{calc}$ , of the number of dodecahedral Voronoi polyhedra per unit cell, deduced by the above authors from the values of  $<\Theta>$ , is given in col.10; and equivalent data (Rivier & Aste, 1996) for 3 alloys discovered after the publication of their results are added to the table, in italics, at positions 6, 7, and 8.

The table shows, first, that for each of its 20 values of  $\langle CN \rangle$  there is a value of  $2(1 + n_0/n_3)$  in Table 5A to equal it; and that, if such pairs of numbers are placed in alignment,  $\langle \Theta \rangle$  is seen to equal the corresponding value of  $60(1 + n_3/n_0)$ ;

N, to equal a small integral multiple of the corresponding value of  $n_3$ ; and the nearest whole number to  $x_{calc}$ , to equal the corresponding value of  $29n_3-5n_0$ ;

secondly that, when the entries of the tables are aligned in that way, gaps appear in those of Table 5B, that is, there are entries in Table 5A to which there is nothing to correspond in Table 5B;

and thirdly that, if to these alloys are added the above three subsequently identified ones, they are found to fill gaps in Table 5B, that is, the same kind of correspondence is found to exist between their data and numbers in Table 5A as was found for the alloys of Shoemaker & Shoemaker.

This last fact indicates that the most recent enumeration of tetrahedrally close-packed alloys (Rivier & Aste, 1996) may not be complete, and that further compounds of that type may eventually be found to fill some, if not all of the remaining gaps in Table 5B. Eq.(6) leads to values of N,  $<\Theta>$ , and <CN> and  $_{12}N_3$  (a multiple of the nearest integer to  $x_{calc}$ ) that indicate structures to be expected, of which the most likely to turn up are those with the smallest values of  $n_3$ : Nos.10, 20, 23, and 33 of Tables 5A & B.

Shoemaker & Shoemaker also determined the partition of N into the numbers of constituent 16-, 15-, 14-, and 12-hedra (not shown here); but Table 5A has nothing to correspond to this partition since, as we have seen, its solutions furnish the numbers  $N_0$  and  $N_3$  (and hence  $n_0$  and  $n_3$ ) only; nor has it any numbers to correspond to the crystal's symmetry, or to its lattice constants (Table 5B). These omissions indicate that the proposed algebraic theory is, as yet, powerless to describe some of the most characteristic properties of a crystal.

Nevertheless, since the theory, even as it now stands, enables some crystals to be classified and topological features of their structure to be quantitatively predicted, it cannot be dismissed as a useless numerical exercise. On the contrary, if as a result of this investigation a tetrahedrally close-packed alloy of the transition metals is found to possess one of the predicted structures, there will be all the more reason to seek, as Einstein envisaged (Aboav, 1997) and as is here attempted, a description of Nature grounded on a purely numerical axiomatic framework.

(to be continued)

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